

MST207 Past Paper 2004

Part 1

Q1 Use separation of variables: $\int \frac{dy}{1+y^2} = \int x^2 dx \Rightarrow \tan^{-1}(y) = \frac{x^3}{3} + C \Rightarrow y = \tan\left(\frac{x^3}{3} + C\right)$

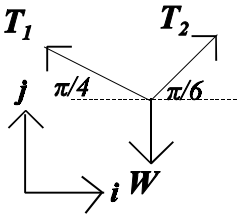
When $x = 0, y = 1$ then $\tan(C) = 1$ and so $C = \pi/4$ So: $y = \tan\left(\frac{x^3}{3} + \frac{\pi}{4}\right)$

Q2 a $\vec{OD} = \frac{\mathbf{0} + (i+0+0) + \left(\frac{i}{2} + \frac{\sqrt{3}j}{2} + 0\right)}{3} = \frac{\left(\frac{3i}{2} + \frac{\sqrt{3}j}{2}\right)}{3} = \left(\frac{1}{2}i + \frac{1}{2\sqrt{3}}j\right)$

b $\vec{OC} = \vec{OD} + \lambda k = \left(\frac{1}{2}i + \frac{1}{2\sqrt{3}}j + \lambda k\right) \Rightarrow |\vec{OC}| = \sqrt{\left(\frac{1}{2}\right)^2 + \left(\frac{1}{2\sqrt{3}}\right)^2 + \lambda^2} = 1$
 $\left(\frac{1}{2}\right)^2 + \left(\frac{1}{2\sqrt{3}}\right)^2 + \lambda^2 = 1 \Rightarrow \lambda = \sqrt{\frac{2}{3}} \Rightarrow |\vec{OC}| = \left(\frac{1}{2}i + \frac{1}{2\sqrt{3}}j + \sqrt{\frac{2}{3}}k\right)$

c $\cos(\theta) = \frac{\vec{OC} \cdot \vec{OD}}{|\vec{OC}| |\vec{OD}|} = \frac{\frac{1}{3}}{\frac{1}{\sqrt{3}}} = \frac{1}{\sqrt{3}}$

Q3



$$\mathbf{W} = -|\mathbf{W}|\mathbf{j} = -mg\mathbf{j}$$

$$\mathbf{T}_1 = -|\mathbf{T}_1|\cos(\pi/4)\mathbf{i} + |\mathbf{T}_1|\sin(\pi/4)\mathbf{j} = -(|\mathbf{T}_1|/\sqrt{2})\mathbf{i} + (|\mathbf{T}_1|/\sqrt{2})\mathbf{j}$$

$$\mathbf{T}_2 = |\mathbf{T}_2|\cos(\pi/6)\mathbf{i} + |\mathbf{T}_2|\sin(\pi/6)\mathbf{j} = (\sqrt{3}|\mathbf{T}_2|/2)\mathbf{i} + (|\mathbf{T}_2|/2)\mathbf{j}$$

In equilibrium $\mathbf{W} + \mathbf{T}_1 + \mathbf{T}_2 = \mathbf{0}$

$$\text{Resolving in } \mathbf{j}\text{-direction: } -mg\mathbf{j} + (|\mathbf{T}_1|/\sqrt{2})\mathbf{j} + (|\mathbf{T}_2|/2)\mathbf{j} = \mathbf{0}$$

$$\text{So: } mg = |\mathbf{T}_1|/\sqrt{2} + |\mathbf{T}_2|/2$$

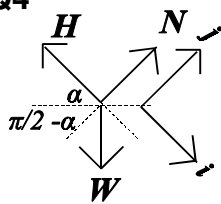
$$\text{Resolving in } \mathbf{i}\text{-direction: } -|\mathbf{T}_1|/\sqrt{2}\mathbf{i} = g\mathbf{i} \text{ so } |\mathbf{T}_1| = \sqrt{2}g; \quad -(|\mathbf{T}_1|/\sqrt{2})\mathbf{i} + (\sqrt{3}|\mathbf{T}_2|/2)\mathbf{i} = \mathbf{0}$$

$$\text{So: } |\mathbf{T}_1|/\sqrt{2} = \sqrt{3}|\mathbf{T}_2|/2 \text{ and so } |\mathbf{T}_1| = \sqrt{3}|\mathbf{T}_2|/\sqrt{2}$$

$$\text{Substituting: } mg = \sqrt{3}|\mathbf{T}_2|/(\sqrt{2})^2 + |\mathbf{T}_2|/2 = \frac{1}{2}(1 + \sqrt{3})|\mathbf{T}_2| \text{ So: } |\mathbf{T}_2| = 2mg/(1 + \sqrt{3}) \text{ as required.}$$

$$\text{Substituting this gives: } |\mathbf{T}_1| = (\sqrt{6}mg)/(1 + \sqrt{3}) \text{ as required.}$$

Q4



$$\mathbf{H} = -k(x - l_0)\hat{\mathbf{s}} = -k(x - l_0)\mathbf{i} \text{ since } \hat{\mathbf{s}} = \mathbf{i}$$

$$\mathbf{N} = |\mathbf{N}|\mathbf{j}$$

$$\mathbf{W} = |\mathbf{W}|\sin(\alpha)\mathbf{i} - |\mathbf{W}|\cos(\alpha)\mathbf{j} = mg \sin(\alpha)\mathbf{i} - mg \cos(\alpha)\mathbf{j}$$

$$\ddot{\mathbf{x}} = \ddot{x}\mathbf{i}$$

$$\text{Since we have motion then: } m\ddot{\mathbf{x}} = \mathbf{H} + \mathbf{N} + \mathbf{W}$$

$$\text{Resolving in } \mathbf{j}\text{-direction gives: } 0 = |\mathbf{N}|\mathbf{j} - |\mathbf{W}|\cos(\alpha)\mathbf{j}$$

$$\text{So: } |\mathbf{N}| = |\mathbf{W}|\cos(\alpha) = mg \cos(\alpha)$$

$$\text{Resolving in } \mathbf{i}\text{-direction gives: } m\ddot{x}\mathbf{i} = -k(x - l_0)\mathbf{i} + mg \sin(\alpha)\mathbf{i}$$

$$m\ddot{x} = -k(x - l_0) + mg \sin(\alpha)$$

$$\text{So equation of motion is: } \ddot{x} + kx/m = g \sin(\alpha) + kl_0/m$$

Equilibrium position, $x = l_{eq}$, is when mass is stationary; i.e. when $\ddot{x} = 0$. Substituting gives:

$$l_{eq} = (mg \sin(\alpha) + kl_0)/k$$

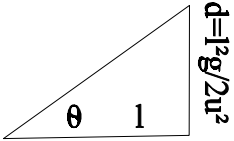
Q5 a Initially $\theta = 0$. Since the trajectory is: $y = h + x \tan \theta - \frac{gx^2}{2u^2}(1 + \tan^2 \theta)$

Simplifying gives:

$$y = h + 0 - \frac{l^2 g}{2u^2}(1 + 0) \text{ or } y = h - \frac{l^2 g}{2u^2}$$

Since the target is at the same height as the launch then vertical distance that the shot misses the target is: $d = \frac{l^2 g}{2u^2}$

b This time the shot is aimed at a height d above the target. So the tangent of the angle of projection is: $\tan\theta = \frac{lg}{2u^2}$ Substituting in the trajectory equation gives:

$$y = h + lx \frac{lg}{2u^2} - \frac{gl^2}{2u^2} \left(1 + \left(\frac{lg}{2u^2} \right)^2 \right) \Rightarrow y = h - \frac{l^4 g^3}{8u^6} \quad \text{and} \quad \frac{l^4 g^3}{8u^6} \neq 0$$


It will only hit the target if the trajectory, y , equals h at l . Since the expression is not zero then the object will fall below the target. We know the expression is not zero since it is composed of all positive elements.

Q6 $A = 8\text{m}^2$; $\Theta_{in} = 19^\circ\text{C}$; $\Theta_{out} = 4^\circ\text{C}$; $b = 0.1\text{m}$; $k = 0.6\text{Wm}^{-1}\text{K}^{-1}$; $h_{in} = 10\text{Wm}^{-2}\text{K}^{-1}$
 $h_{out} = 150\text{Wm}^{-2}\text{K}^{-1}$

a $q = \left(\frac{1}{h_{in}} + \frac{1}{h_{out}} + \frac{b}{k} \right)^{-1} A(\Theta_{in} - \Theta_{out}) \Rightarrow q \approx 3.66 \times 8 \times 15 \approx 439.02\text{W}$

b $q = h_{in} A(\Theta_{in} - \Theta_1) \Rightarrow \Theta_1 = \Theta_{in} - \frac{q}{Ah_{in}} = 19 - \frac{439.02}{10 \times 8} \approx 13.5^\circ\text{C}$

Q7

$$\begin{array}{l} \mathbf{R}_1 \left[\begin{array}{ccc|c} 1 & 2 & 1 & 2 \\ 2 & 2 & 1 & a \\ 2 & 0 & b & 4 \end{array} \right] \quad \mathbf{R}_1 \left[\begin{array}{ccc|c} 1 & 2 & 1 & 2 \\ 0 & -2 & -1 & a-4 \\ 0 & -4 & b-2 & 0 \end{array} \right] \quad \mathbf{R}_1 \left[\begin{array}{ccc|c} 1 & 2 & 1 & 2 \\ 0 & -2 & -1 & a-4 \\ 0 & 0 & b & -2(a-4) \end{array} \right] \\ \mathbf{R}_2 \left[\begin{array}{ccc|c} 1 & 2 & 1 & 2 \\ 2 & 2 & 1 & a \\ 2 & 0 & b & 4 \end{array} \right] \quad \mathbf{R}_{2a} = \mathbf{R}_2 - 2\mathbf{R}_1 \quad \mathbf{R}_{2a} \left[\begin{array}{ccc|c} 1 & 2 & 1 & 2 \\ 0 & -2 & -1 & a-4 \\ 0 & 0 & b & -2(a-4) \end{array} \right] \\ \mathbf{R}_3 \left[\begin{array}{ccc|c} 1 & 2 & 1 & 2 \\ 2 & 2 & 1 & a \\ 2 & 0 & b & 4 \end{array} \right] \quad \mathbf{R}_{3a} = \mathbf{R}_3 - 2\mathbf{R}_1 \quad \mathbf{R}_{3b} = \mathbf{R}_{3a} - 2\mathbf{R}_{2a} \end{array}$$

For there to be infinitely many solutions then all the entries on row \mathbf{R}_{3b} have to be zero.
Hence: $b = 0$ and: $-2(a - 4) = 0$. So $2a = 8$ So: $a = 4$

Q8 a So corresponding eigenvalue is 12.

$$\begin{bmatrix} 10 & 2 & 0 \\ 4 & 11 & 6 \\ -2 & -1 & 6 \end{bmatrix} \begin{bmatrix} 2 \\ 2 \\ -1 \end{bmatrix} = \begin{bmatrix} 24 \\ 24 \\ -12 \end{bmatrix} = 12 \begin{bmatrix} 2 \\ 2 \\ -1 \end{bmatrix}$$

b

Let $\lambda = 6$; $(10 - 6)x + 2y + 0 = 2$; $4x + (11 - 6)y + 6z = 0$; $-2x - y + (6 - 6)z = 0$
The first and third equations give $-2x = y$. So if $x = 1$ the eigenvector so far is $[1 \ -2 \ z]^T$
Substituting these into equation 2 gives $z = 1$. So corresponding eigen vector is $[1 \ -2 \ 1]^T$

Q9

$$\begin{array}{l} \Rightarrow \mathbf{i} \quad \mathbf{H}_{it} = -k(x - l_0)\mathbf{i} \\ \leftarrow \mathbf{i} \quad \mathbf{H}_{it} = -2k(d - x - l_0)(-\mathbf{i}) = 2k(d - x - l_0)\mathbf{i} \\ \mathbf{H}_{it} \quad \mathbf{H}_{it} \quad \mathbf{R} \quad \mathbf{R} = -r\mathbf{x}\mathbf{i} \end{array}$$

$m\ddot{x}\mathbf{i} = -k(x - l_0)\mathbf{i} + 2k(d - x - l_0)\mathbf{i} - r\mathbf{x}\mathbf{i}$ which results in: $m\ddot{x} + r\dot{x} + 3kx = k(2d - l_0)$, eqn of motion.

Q10 a By principle of Conservation of momentum: $m(\mathbf{u}\mathbf{i} + \mathbf{v}\mathbf{j}) + m\mathbf{0} = m\mathbf{p}\mathbf{i} + m\mathbf{q}\mathbf{j}$. Resolving into \mathbf{i} - and \mathbf{j} -directions gives $p = u$ and $q = v$ and so after collision one particle has velocity $u\mathbf{i}$ and the other $v\mathbf{j}$.

b If elastic then K.E. before = K.E. after. K.E. before = $\frac{1}{2}mu^2$ where $u = |\mathbf{u}| = \sqrt{(u^2 + v^2)}$ so K.E. before = $\frac{1}{2}m(u^2 + v^2)$.

K.E. after = $\frac{1}{2}mp^2 + \frac{1}{2}mq^2$ where $p = |\mathbf{p}| = u$ and $q = |\mathbf{q}| = v$ so K.E. after = $\frac{1}{2}m(u^2 + v^2)$. Hence collision is elastic as K.E. before = K.E. after.

Q11 a $f_x = 8xy + y^2 + 12$ $f_y = 4x^2 + 2xy = x(4x + 2y)$

Stationary points are found by solving the simultaneous equations $f_x = 0$ and $f_y = 0$. i.e.

$8xy + y^2 + 12 = 0$ and $x(4x + 2y) = 0$

The second one gives a solution of $x = 0$. Substituting this into first equation gives $y^2 = -12$ which has solutions $\pm 2\sqrt{3}i$ and so $(0, +2\sqrt{3}i)$ and $(0, -2\sqrt{3}i)$ are two stationary points.

Now: $(4x + 2y) = 0$ gives: $-2x = y$. Substituting into $f_x = 0$ gives $x = \pm 1$. When $x = +1$, $y = -2$ and when $x = -1$, $y = +2$ (found by substituting values of x into $f_y = 0$)

So all stationary points are: $(0, +2\sqrt{3}i)$; $(0, -2\sqrt{3}i)$; $(1, -2)$; $(-1, 2)$

b $A = f_{xx} = 8y$ $C = f_{yy} = 2x$ $B = f_{xy} = 8x + 2y$

With point where x is positive: $A = f_{xx}(1, -2) = -16$; $C = f_{yy}(1, -2) = 2$; $B = f_{xy}(1, -2) = 8 - 4 = 4$

$AC - B^2$ at $(1, -2) = -32 + 16 = -16 < 0$ hence stationary point is a Saddle Point.

Q12 i Plot Y_N against h^4

ii $C \approx (Y_2 - Y_1)/(h_1^4 - h_2^4) = (0.370098 - 0.368018)/((0.25)^4 - (0.125)^4) \approx 0.568$

For 15 d.p. accuracy: $h^4 < 0.5 \times 10^{-15}/C$ Thus $h < 0.000172$

Q13 To find the area we let f , the surface density function, = 1. So:

$$A = \int_{x=0}^{x=\pi/2} \int_{y=1-2x/\pi}^{y=\cos x} 1 \, dy \, dx + \int_{x=\pi/2}^{x=\pi} \int_{y=\cos x}^{y=1-2x/\pi} 1 \, dy \, dx = \int_{x=0}^{x=\pi/2} \left(\cos x - 1 - \frac{2x}{\pi} \right) dx + \int_{x=\pi/2}^{x=\pi} \left(1 - \frac{2x}{\pi} - \cos x \right) dx$$

$$= \left[\sin x - x - \frac{x^2}{\pi} \right]_0^{\pi/2} + \left[x - \frac{x^2}{\pi} - \sin x \right]_{\pi/2}^{\pi} = 2 - \frac{\pi}{2}$$

Q14 Since rotating about fixed axis then $K.E = \frac{1}{2}I\omega^2 = 100 \text{ J}$

$I = I_{rod} + 2I_{disk} = \frac{1}{2}m_{rod}r_{rod}^2 + M_{disk}R_{disk}^2$

So:

$$\omega^2 = \frac{100}{\frac{1}{2}(1)(0.02)^2 + 3(0.25)^2} \approx 532.765$$

$$\omega = \sqrt{\omega^2} \approx 23.08 \text{ rads}^{-1}$$

Part 2

Q15 i Characteristic equations is: $4\lambda^2 + 4\lambda + 5 = 0$ which has solutions $-\frac{1}{2} \pm i$ and so homogeneous solution is: $x = e^{-\frac{1}{2}t}(C \cos(t) + D \sin(t))$

Let: $x_p = ke^{-\frac{1}{2}t}$; then: $x' = -\frac{1}{2}ke^{-\frac{1}{2}t}$ and: $x'' = \frac{1}{4}ke^{-\frac{1}{2}t}$

Substituting gives $k = 1$, so solutions so far is: $x = e^{-\frac{1}{2}t}(C \cos(t) + D \sin(t)) + e^{-\frac{1}{2}t}$

Let: $x_p = at^2 + bt + c$; then: $x' = 2at + b$; and: $x'' = 2a$

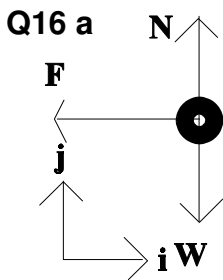
Substituting gives $a = 1$, $b = 0$ and $c = 1$ so solutions so far is:

$x = e^{-\frac{1}{2}t}(C \cos(t) + D \sin(t)) + e^{-\frac{1}{2}t} + t^2 + 1$

Since $x(0) = 0$, substituting gives: $0 = C + 1 + 1$ so $C = -2$. So the solution so far is:

$x = e^{-\frac{1}{2}t}(D \sin(t) - 2\cos(t)) + e^{-\frac{1}{2}t} + t^2 + 1$

Since $x'(0) = 0$ then $x' = e^{-\frac{1}{2}t}(D\cos(t) + 2\sin(t) - \frac{1}{2}(D\sin(t) - 2\cos(t) + 1)) + 2t$ which, on substitution gives $D = -\frac{1}{2}$. So the particular solution is: $x = e^{-\frac{1}{2}t}(1 - 2\cos(t) - \frac{1}{2}\sin(t)) + t^2 + 1$



$\ddot{\mathbf{x}} = \ddot{x}\mathbf{i}$; $\mathbf{v} = v\mathbf{i}$; $\mathbf{N} = |\mathbf{N}|\mathbf{j}$; $\mathbf{F} = -|\mathbf{F}|\mathbf{i}$; $|\mathbf{F}| = \mu|\mathbf{N}|$; $\mathbf{W} = -|\mathbf{W}|\mathbf{j} = -mg\mathbf{j}$

Since there is motion we have: $\mathbf{N} + \mathbf{F} + \mathbf{W} = m\ddot{\mathbf{x}}$

Resolving in \mathbf{j} -direction gives: $|\mathbf{N}| = mg$. Consequently $|\mathbf{F}| = \mu mg$

Since we have motion then by Newton's Second Law, resolving in \mathbf{i} -direction gives:

$m\ddot{x}\mathbf{i} = -|\mathbf{F}|\mathbf{i} = -\mu mg\mathbf{i}$. This simplifies to: $\ddot{x} = -\mu g$

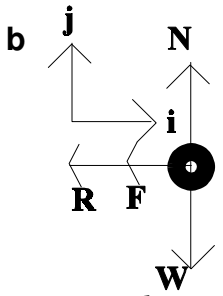
$$\ddot{x} = v \frac{dv}{dx} = -\mu g \Rightarrow \int v dv = -\mu g \int dx \Rightarrow \frac{v^2}{2} = -\mu gx + C$$

Initially when $x = 0$, $v = u_0$ So:

$$\frac{u_0^2}{2} = C \Rightarrow \frac{v^2}{2} = \frac{u_0^2}{2} - \mu gx \Rightarrow v^2 = u_0^2 - 2\mu gx \Rightarrow v = \sqrt{u_0^2 - 2\mu gx} \quad \text{or: } x = \frac{u_0^2 - v^2}{2\mu g}$$

When has come to rest, $v = 0$. Substituting this, and all the other given values for the constants gives x as:

$$x = \frac{u_0^2 - v^2}{2\mu g} = \frac{2^2 - 0}{2 \times 0.02 \times 9.81} \approx 10.19\text{m (to 2.d.p.)}$$



$\mathbf{R} = -|\mathbf{R}|\mathbf{i} = -kv^2\mathbf{i}$

All other forces and vectors defined as above and $|\mathbf{F}| = \mu mg$ as before.

Since there is motion we have: $\mathbf{N} + \mathbf{F} + \mathbf{R} + \mathbf{W} = m\ddot{\mathbf{x}}$

Resolving in \mathbf{i} -direction now gives: $m\ddot{x}\mathbf{i} = -|\mathbf{F}|\mathbf{i} - |\mathbf{R}|\mathbf{i} = -\mu mg\mathbf{i} - kv^2\mathbf{i}$

Hence: $\ddot{x} = -\mu g - kv^2/m$

$k = c_2 D^2 = 0.2 \times (0.05)^2 = 0.0005$ and $k/m = 0.0005/2 = 0.00025 = 1/4000$

$$\ddot{x} = v \frac{dv}{dx} = -\mu g - \frac{v^2}{4000} \Rightarrow \int \frac{v dv}{\mu g + \frac{v^2}{4000}} = 2000 \int \frac{\frac{v dv}{2000}}{\frac{4000\mu g + v^2}{4000}} = 2000 \int \frac{2v dv}{4000\mu g + v^2} = -\int dx$$

$$2000 \ln(4000\mu g + v^2) = -x + C \quad (\text{integrate by substitution})$$

Initially, $x = 0$ and $v = u_0$. So:

$$c = 2000 \ln(4000\mu g + u_0^2) \text{ so } x = 2000 \left(\ln(4000\mu g + u_0^2) - \ln(4000\mu g + v^2) \right)$$

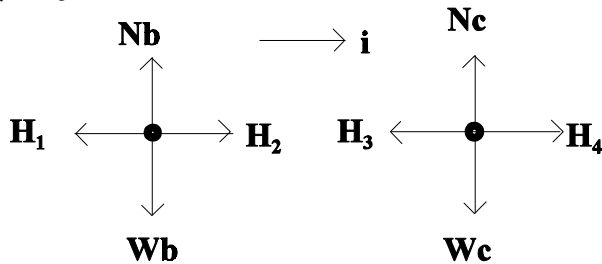
$$x = 2000 \ln \left(\frac{4000\mu g + u_0^2}{4000\mu g + v^2} \right)$$

When comes to rest $v = 0$. Substituting this and all the other given constants gives:

$$x = 2000 \ln \left(\frac{4000 \times 0.02 \times 9.81 + 2^2}{4000 \times 0.02 \times 9.81 + 0} \right) = 2000 \ln(1.0051) \approx 10.17\text{m (to 2.d.p.)}$$

c Adding air resistance stops the mass about 2cm sooner. Two cm over approximately 10.2m represents about 0.2%. So the effect of air resistance accounts for about 0.2% of the resistive force and friction for 99.8% of the resistive force. Air resistance is effectively negligible in this case.

Q17 a



$$\begin{aligned} \Delta H_1 &= -kxi \\ \Delta H_2 &= 4k(y - x)i \\ \Delta H_3 &= -\Delta H_2 \\ \Delta H_4 &= -kyi \\ m\ddot{x}i &= \Delta H_1 + \Delta H_2 = -5kxi + 4kyi \\ \text{So: } \ddot{x} &= k/m(-5x + 4y) \\ \text{Sim: } \ddot{y} &= k/m(4x - 5y) \end{aligned}$$

So:
$$\begin{bmatrix} \ddot{x} \\ \ddot{y} \end{bmatrix} = \frac{k}{m} \begin{bmatrix} -5 & 4 \\ 4 & -5 \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$$

b Eigenvalues of the matrix are -1 and -9 and hence of the whole system are $-k/m$ and $-9k/m$. The normal mode angular frequencies are given by $\sqrt{-\lambda}$, where λ is an eigenvalue of the system. So the normal mode angular frequencies are $\sqrt{k/m}$ and $3\sqrt{k/m}$

c Eigenvectors are $[1, 1]$ and $[-1, 1]$ respectively. Hence general solution is:

$$\begin{bmatrix} x(t) \\ y(t) \end{bmatrix} = C_1 \begin{bmatrix} 1 \\ 1 \end{bmatrix} \cos\left(\sqrt{\frac{k}{m}}t + \phi\right) + C_2 \begin{bmatrix} -1 \\ 1 \end{bmatrix} \cos\left(3\sqrt{\frac{k}{m}}t + \phi\right)$$

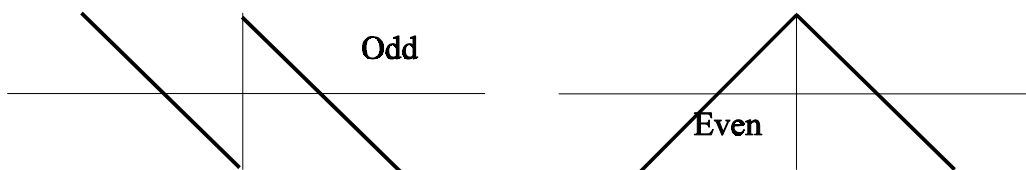
d All phase angles, ϕ , are zero. At $t = 0$, then $x(0) = 0.1\text{m}$ and $y(0) = 0.2\text{m}$. Substituting gives $C_1 = 0.15\text{m}$ (15 cm) and $C_2 = 0.05\text{m}$ (5cm). Hence at time t , the positions $x(t)$ and $y(t)$ are given (in metres) by:

$$\begin{bmatrix} x(t) \\ y(t) \end{bmatrix} = 0.15 \begin{bmatrix} 1 \\ 1 \end{bmatrix} \cos\left(\sqrt{\frac{k}{m}}t\right) + 0.05 \begin{bmatrix} -1 \\ 1 \end{bmatrix} \cos\left(3\sqrt{\frac{k}{m}}t\right)$$

Q18. This is virtually identical to the 2001 paper Q18

Q 19 a

$$f_{\text{odd}}(t) = \begin{cases} 1/2 - t & 0 \leq t \leq 1 \\ -(1/2 - t) & -1 \leq t \leq 0 \end{cases} \quad f_{\text{even}}(t) = \begin{cases} 1/2 - t & 0 \leq t \leq 1 \\ 1/2 + t & -1 \leq t \leq 0 \end{cases}$$



b

$$A_0 = \int_0^1 (1/2 - t) dt = [1/2t - 1/2t^2]_0^1 = 0$$

$$A_r = 2 \int_0^1 (1/2 - t) \cos(r\pi t) dt = \frac{2}{r^2 \pi^2} [-\cos(r\pi t)]_0^1 = \frac{2}{r^2 \pi^2} (1 - \cos(r\pi)) \quad (r = 1, 2, \dots)$$

c

$$B_r = 2 \int_0^1 (1/2 - t) \sin(r\pi t) dt = \frac{1 + \cos(r\pi)}{r\pi} \quad (A_r \text{ and } B_r \text{ by integration by parts})$$

d The cosine function would be better because the terms decrease by a factor of $1/r^2$ compared to the sine function which decreases by a factor of only $1/r$ each time.

Q 20 a Path is a circle centred on the origin with radius = 1

b $\dot{\mathbf{r}}(t) = -\sin(t)\mathbf{i} + \cos(t)\mathbf{j} + 0\mathbf{k}$ and $\mathbf{F}(\mathbf{r}) = \alpha^2\sin(t)\mathbf{i} + \cos(t)\mathbf{j} + \sin(t)\cos(t)\mathbf{k}$

$$\int_C \mathbf{F} \cdot d\mathbf{r} = \int_0^{2\pi} \mathbf{F}(\mathbf{r}) \cdot \dot{\mathbf{r}}(t) dt = \pi(1 - \alpha^2)$$

c If the field is conservative then the line integral of any closed loop is zero. The above is a line integral of a closed loop (circle). This is zero when $\alpha^2 = 1$ hence when $\alpha = \pm 1$. It is not conservative for all other values of α .

d $\text{curl } \mathbf{F} = (1 - \alpha^2)\mathbf{k}$ and so \mathbf{F} is conservative only when $\alpha = \pm 1$.